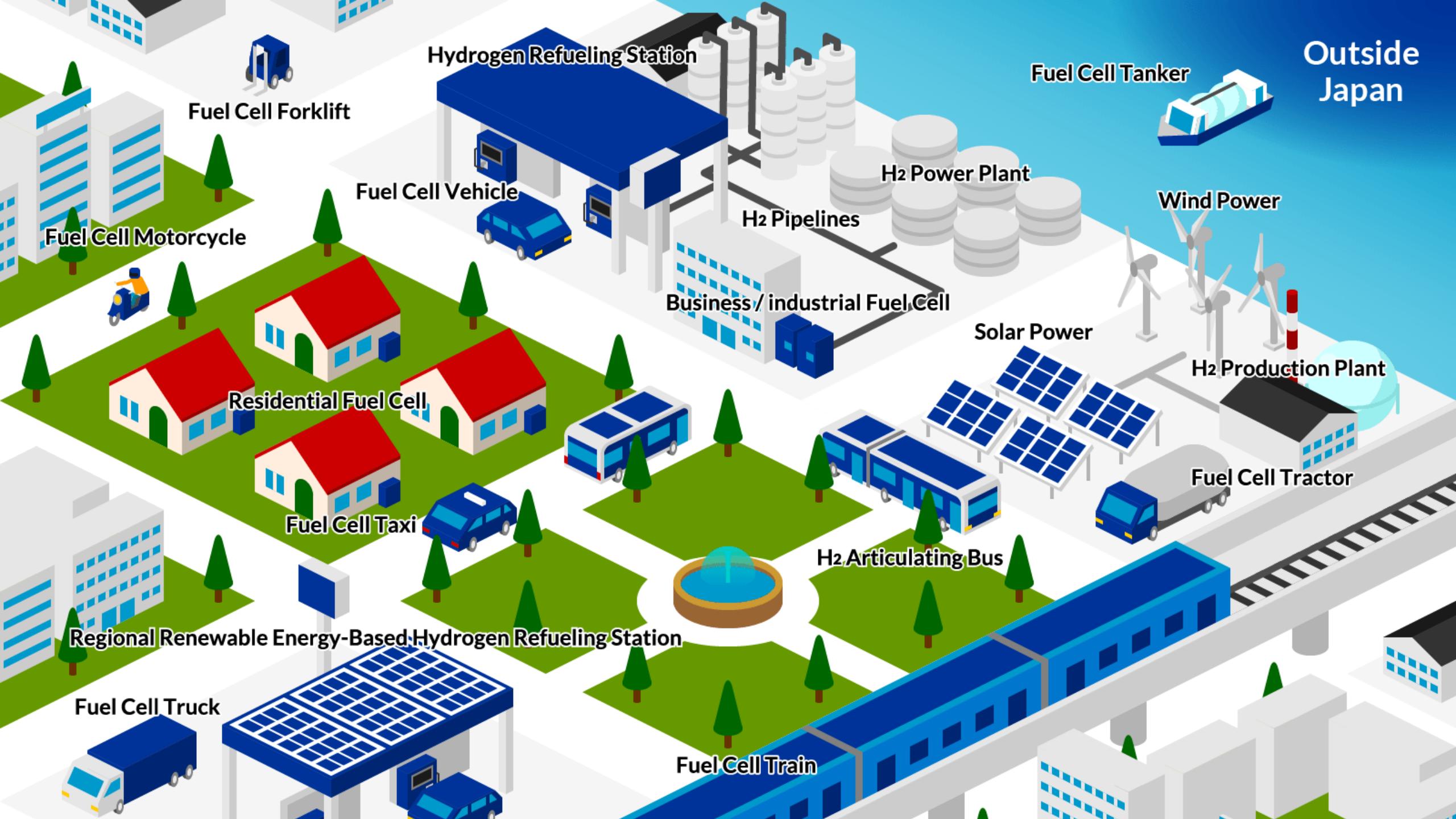


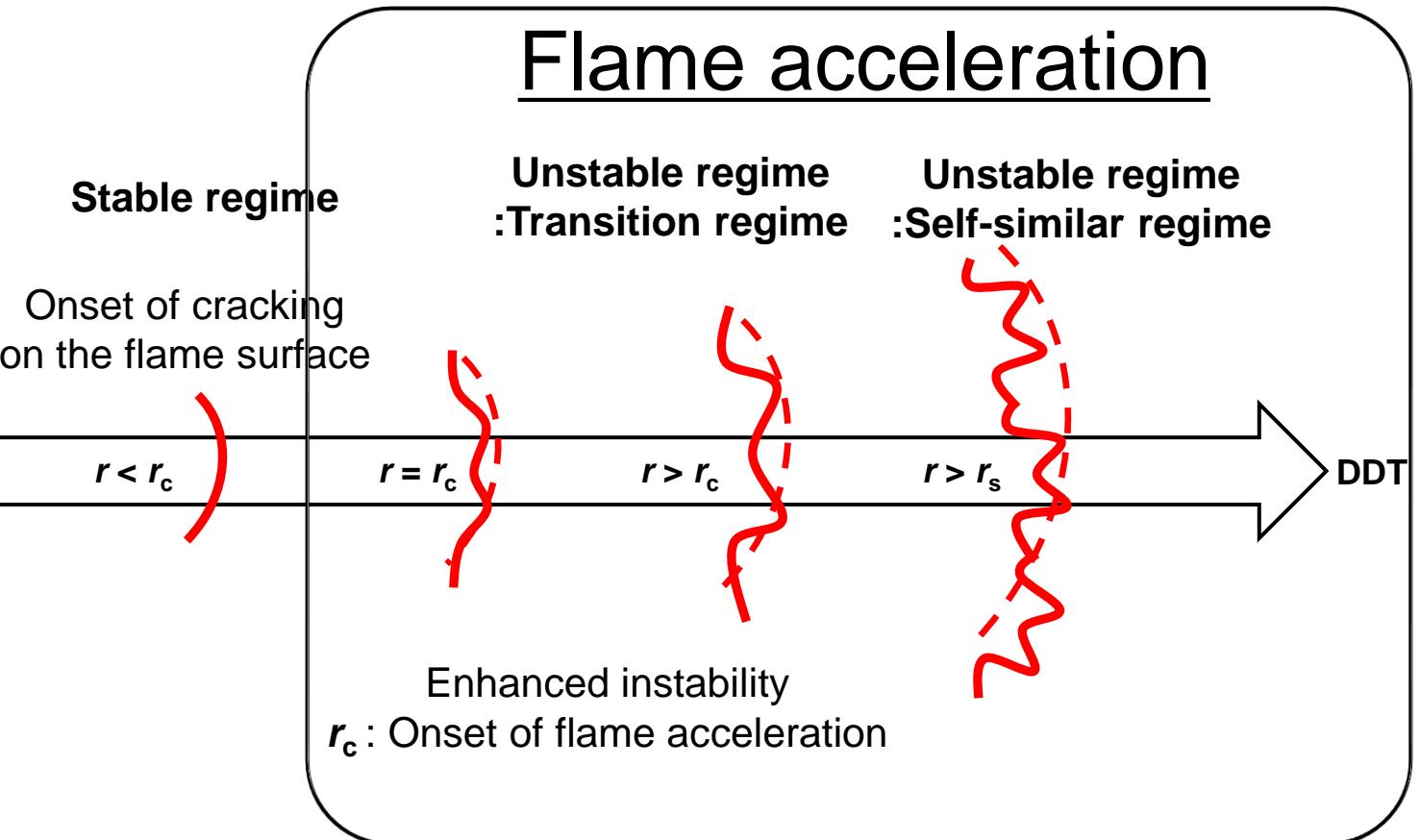
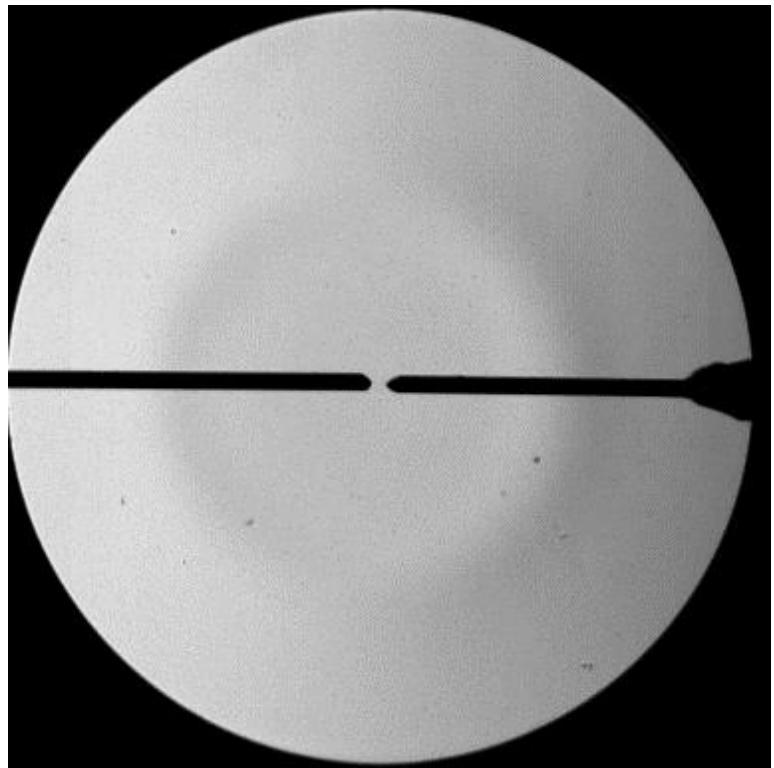
# Flame propagation characteristics in hydrogen-air mixtures

Wookyung Kim

Department of Mechanical Systems Engineering  
Hiroshima University



# How does a flame propagate?



# Flame acceleration

Cellular instabilities

- Darrieus–Landau instability
- Diffusive-thermal instability

Wrinkled flame front

Larger surface area

Flame acceleration

$$r \propto t^\alpha$$

Stable regime

Onset of cracking  
on the flame surface

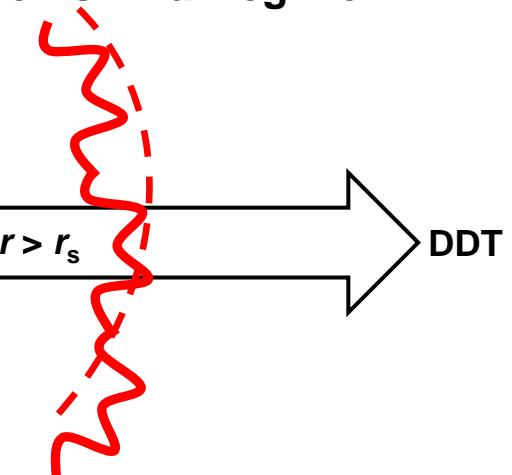
$$r < r_c$$

Unstable regime  
:Transition regime

$$r = r_c$$

$$r > r_c$$

Unstable regime  
:Self-similar regime



Enhanced instability

$r_c$ : Onset of flame acceleration

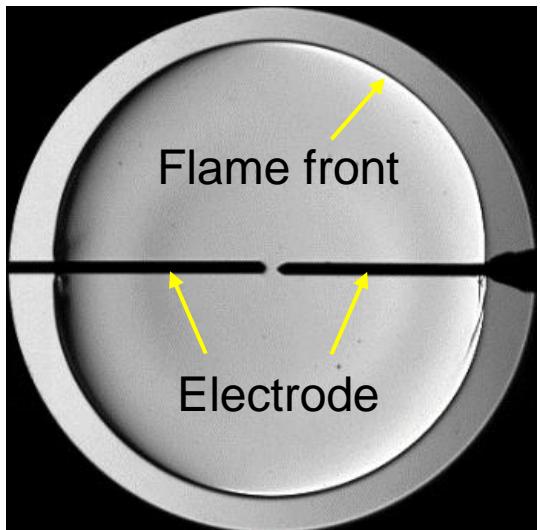
- W. Kim et al., *International Journal of Hydrogen Energy*, 43 (2018) 12556-15564,
- W. Kim et al., *Journal of Loss Prevention in the Process Industries*, 60 (2019) 264-268.
- W. Kim et al., *International Journal of Hydrogen Energy*, 45 (2020) 25608-25614.

# Flame acceleration

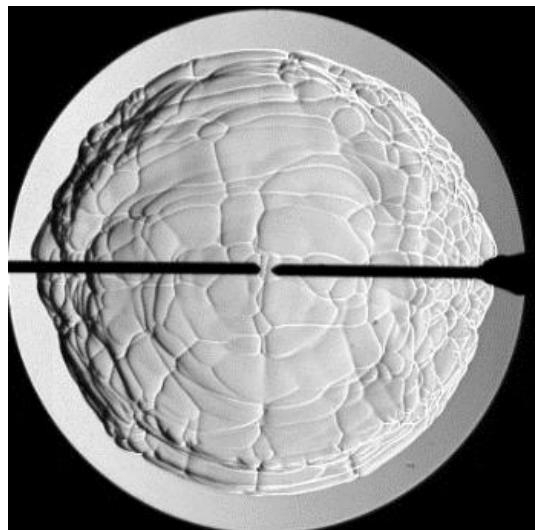
## Cellular instabilities

- Darrieus–Landau instability
- Diffusive-thermal instability

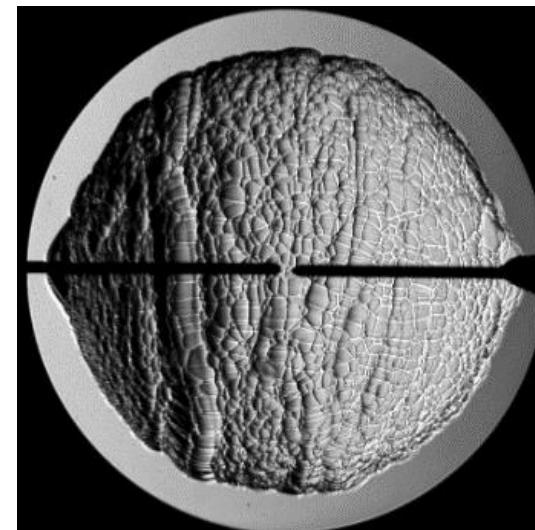
$P_i = 0.1 \text{ MPa}$ ,  $\phi = 2.0$   
 $Le > 1$ ,  $\delta = 0.31 \text{ mm}$



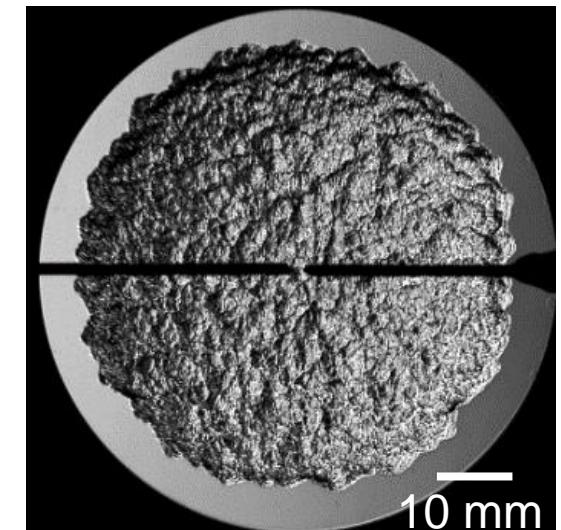
$P_i = 0.1 \text{ MPa}$ ,  $\phi = 0.5$   
 $Le < 1$ ,  $\delta = 0.49 \text{ mm}$



$P_i = 0.5 \text{ MPa}$ ,  $\phi = 2.0$   
 $Le > 1$ ,  $\delta = 0.04 \text{ mm}$



$P_i = 0.5 \text{ MPa}$ ,  $\phi = 0.5$   
 $Le < 1$ ,  $\delta = 0.17 \text{ mm}$



- W. Kim et al., *International Journal of Hydrogen Energy*, 43 (2018) 12556-15564,
- W. Kim et al., *Journal of Loss Prevention in the Process Industries*, 60 (2019) 264-268.
- W. Kim et al., *International Journal of Hydrogen Energy*, 45 (2020) 25608-25614.

# Flame acceleration

Cellular instabilities

- Darrieus–Landau instability
- Diffusive-thermal instability

Wrinkles

Larger

Flame acceleration

$$r \propto t^\alpha$$

- When does self-acceleration occur?  
 $r_c$  : Critical flame radius for onset of flame acceleration
- Does the flame self-accelerate?  
 $r \propto t^\alpha$  i.e.  $\alpha > 1$
- Is self-accelerating flame self-similar?  
 $\alpha = \text{constant}$

Self-similar regime  
Self-similar regime

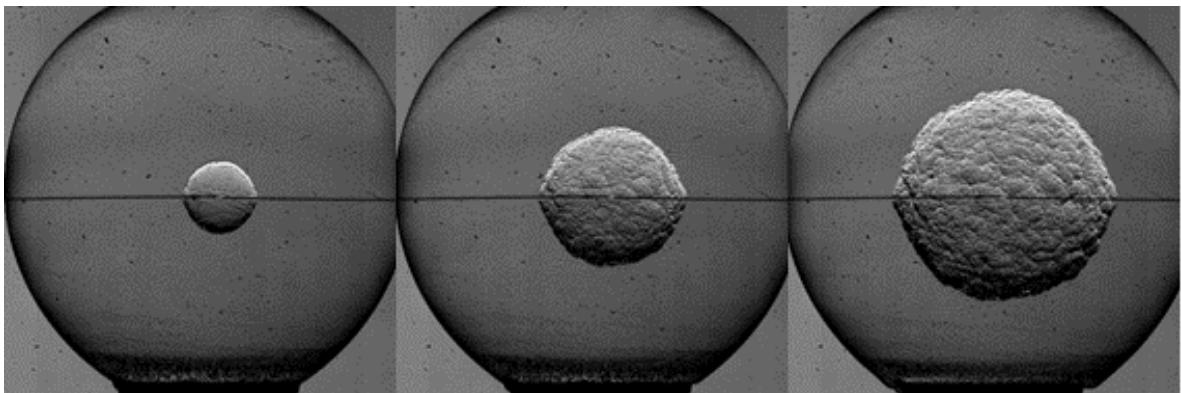
DDT

Enhanced instability  
 $r_c$ : Onset of flame acceleration

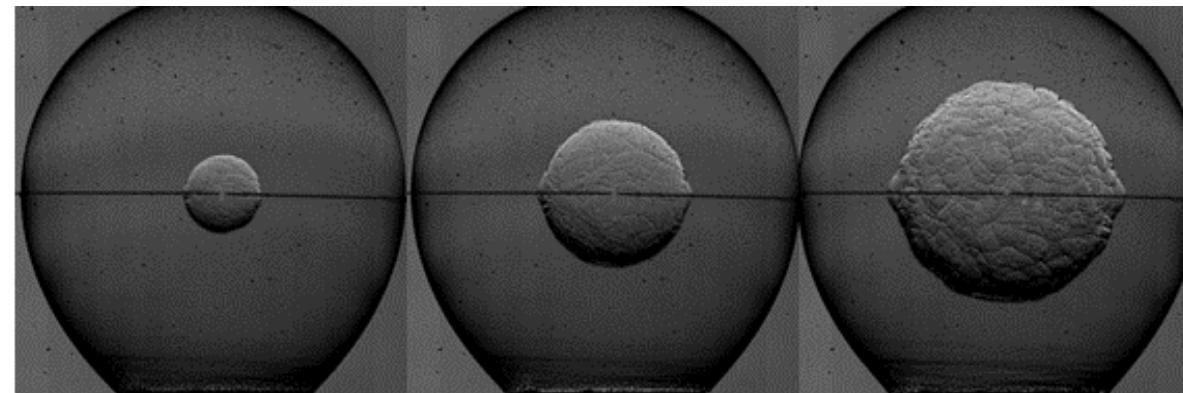


- W. Kim et al., *International Journal of Hydrogen Energy*, 43 (2018) 12556-15564,
- W. Kim et al., *Journal of Loss Prevention in the Process Industries*, 60 (2019) 264-268.
- W. Kim et al., *International Journal of Hydrogen Energy*, 45 (2020) 25608-25614.

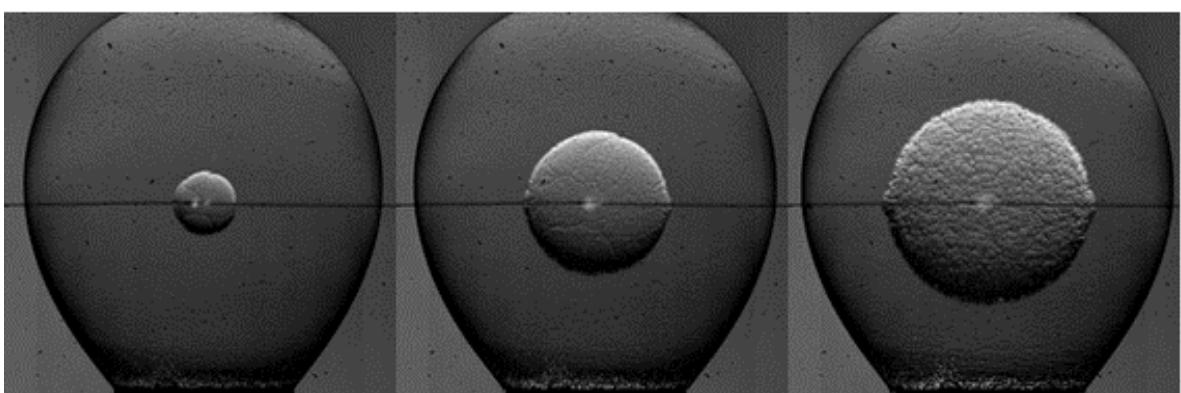
# Hydrogen-oxygen flame



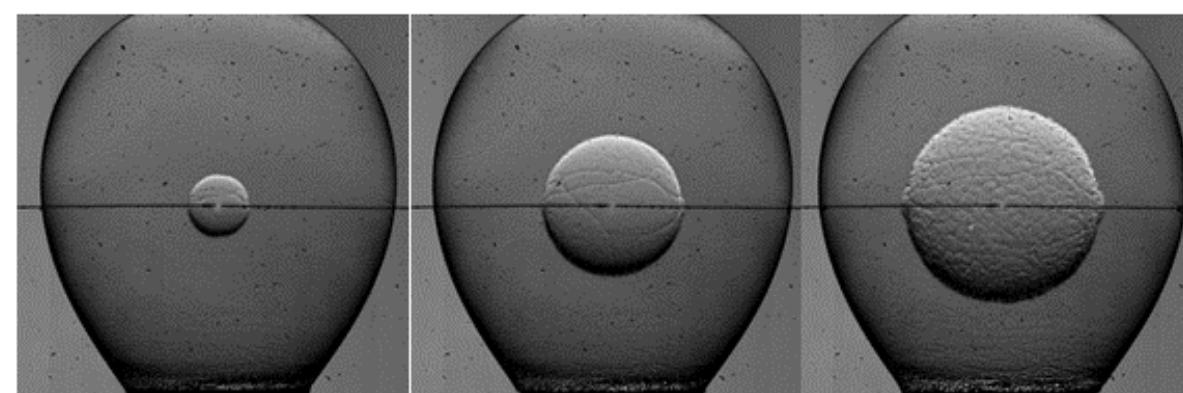
$r = 10.7 \text{ mm}, t = 0.65 \text{ ms}$     $r = 20.1 \text{ mm}, t = 1.20 \text{ ms}$     $r = 30.1 \text{ mm}, t = 1.70 \text{ ms}$   
 $\phi = 0.2$



$r = 10.6 \text{ mm}, t = 0.20 \text{ ms}$     $r = 19.5 \text{ mm}, t = 0.35 \text{ ms}$     $r = 29.0 \text{ mm}, t = 0.50 \text{ ms}$   
 $\phi = 0.6$

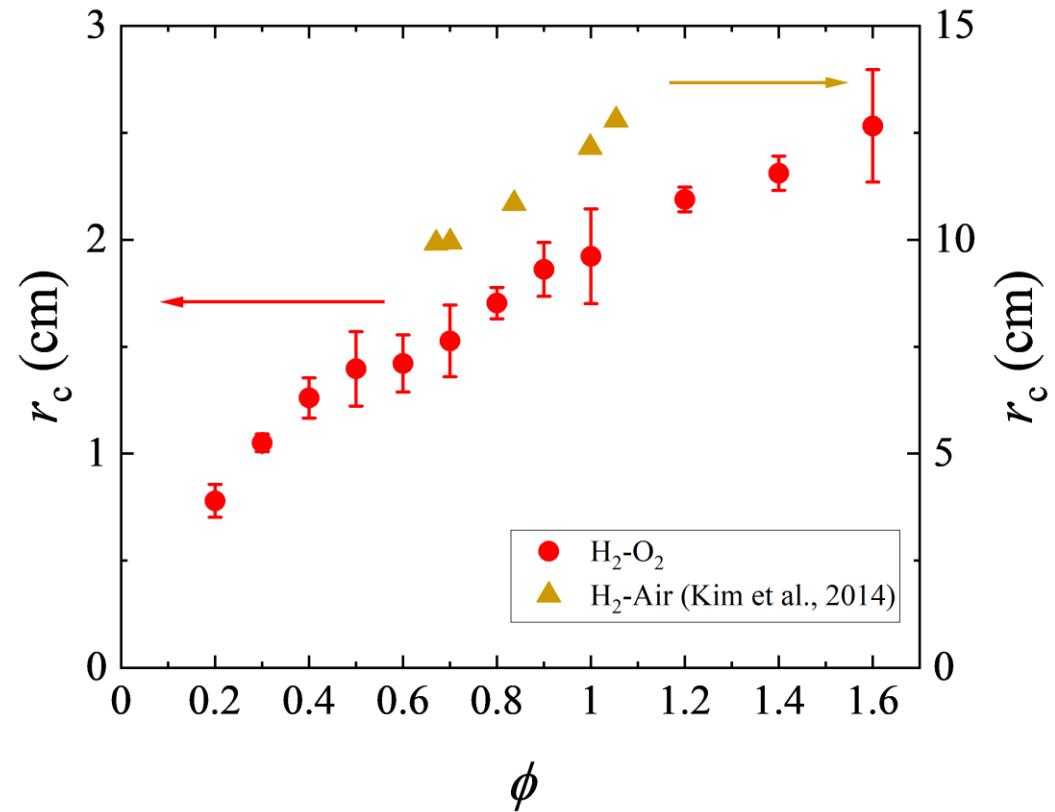
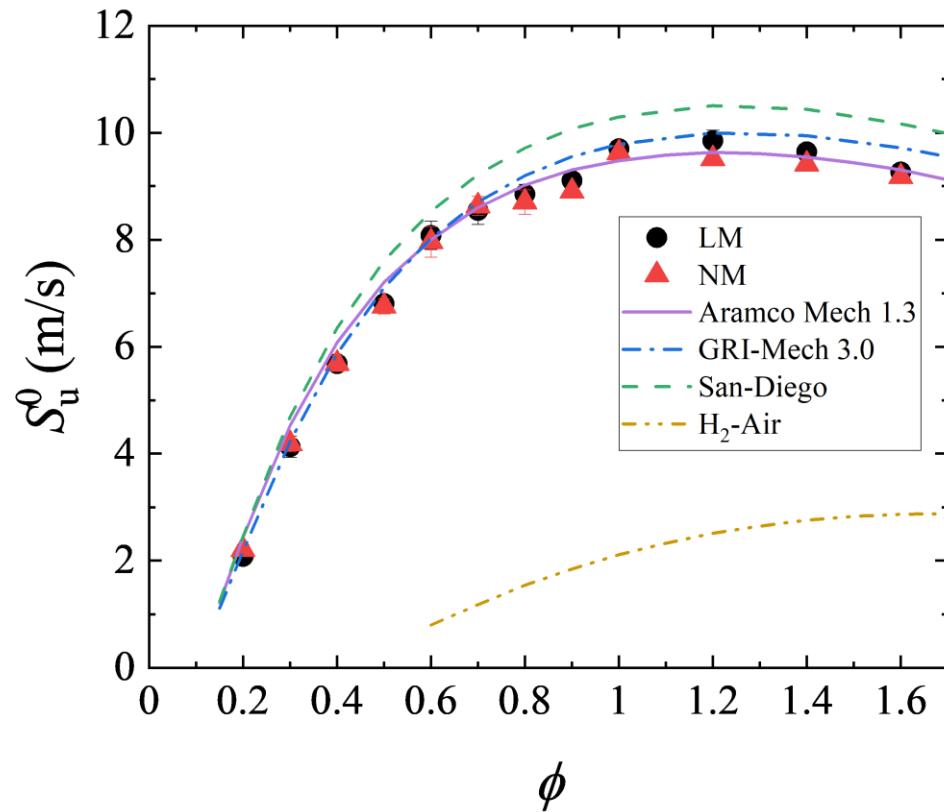


$r = 9.71 \text{ mm}, t = 0.15 \text{ ms}$     $r = 20.8 \text{ mm}, t = 0.30 \text{ ms}$     $r = 29.3 \text{ mm}, t = 0.40 \text{ ms}$   
 $\phi = 1.0$

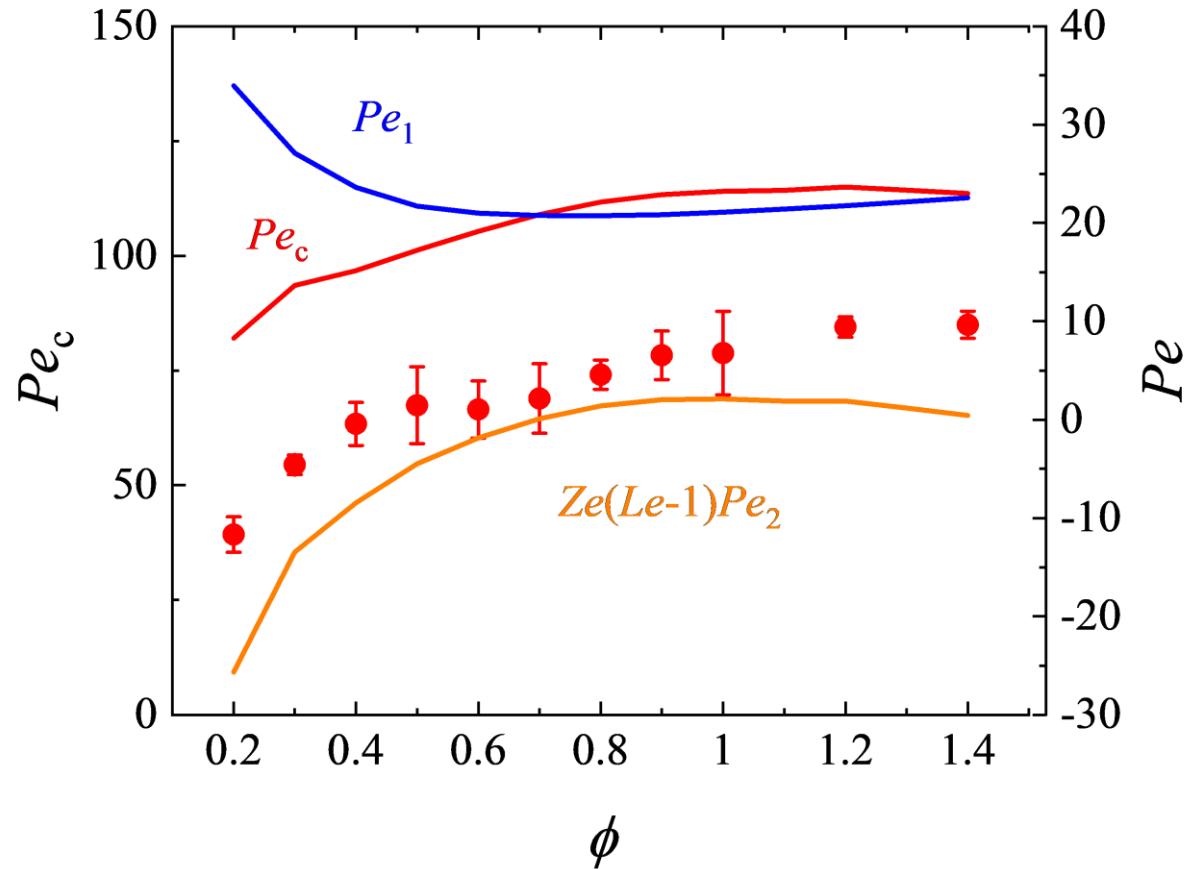


$r = 9.51 \text{ mm}, t = 0.15 \text{ ms}$     $r = 20.3 \text{ mm}, t = 0.30 \text{ ms}$     $r = 28.0 \text{ mm}, t = 0.40 \text{ ms}$   
 $\phi = 1.4$

# Onset of self-acceleration



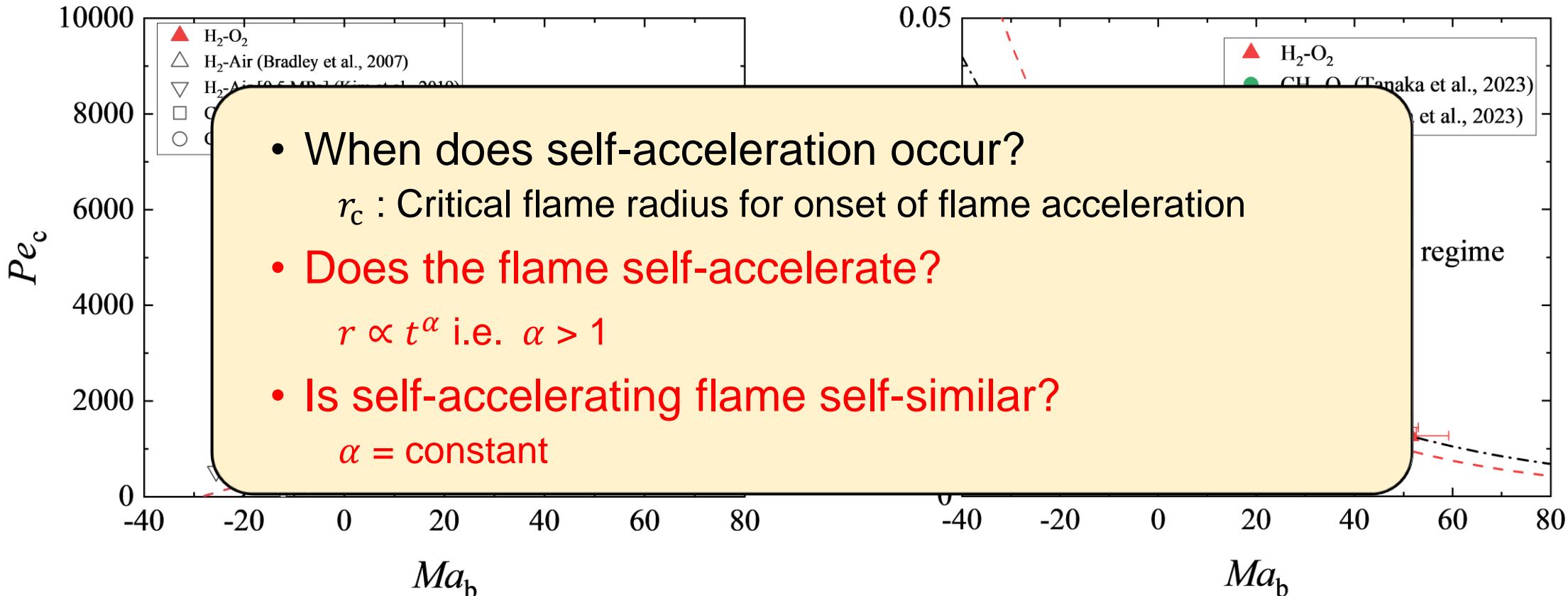
# Onset of self-acceleration



- For the linear analysis of Bechtold and Matalon,  $Pe_c$  expressed the influences due to Darrieus–Landau and diffusive-thermal instabilities.

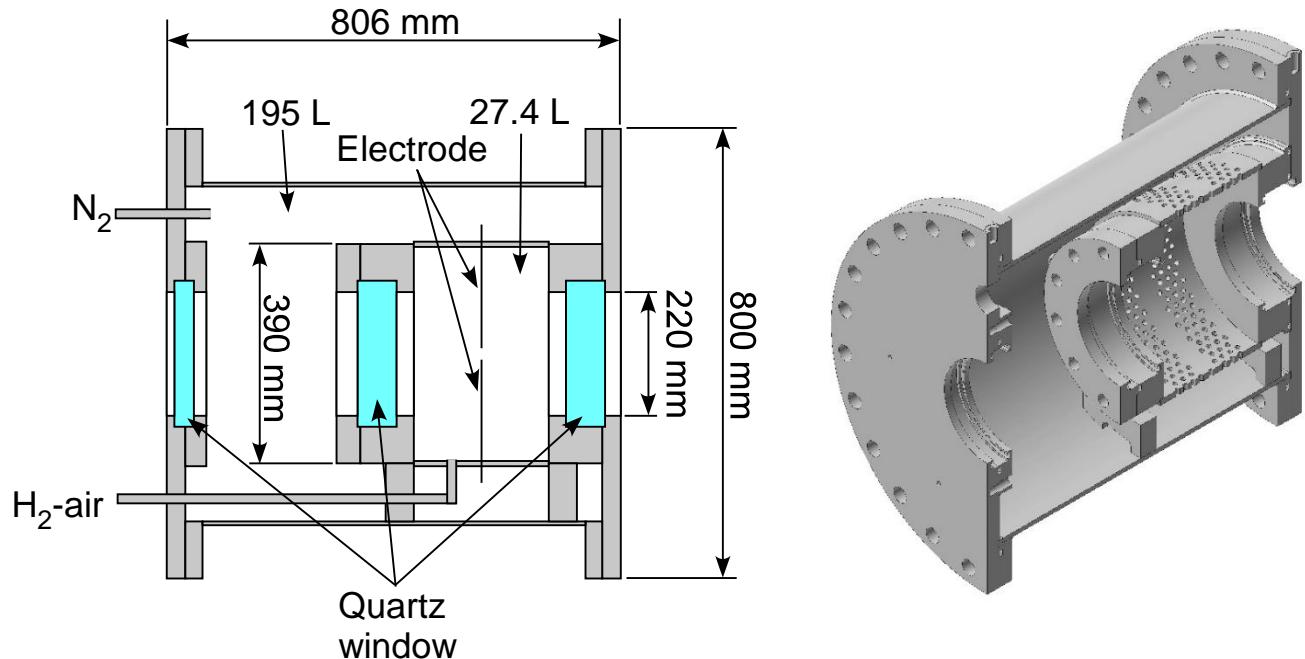
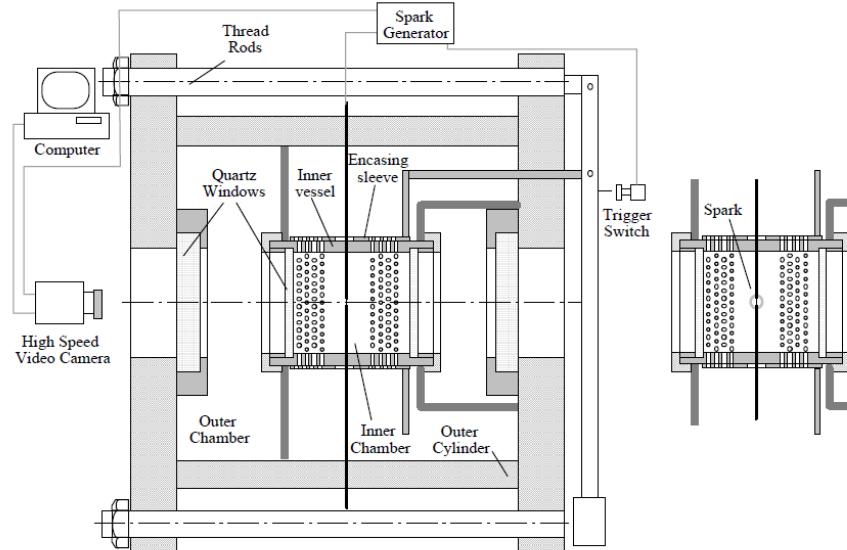
$$Pe_c = Pe_1(\sigma) + Ze(Le - 1)Pe_2(\sigma)$$

# Onset of self-acceleration



- D Bradley et al., Combustion and Flame 149 (2007) 162-172
- W Kim et al., International Journal of Hydrogen Energy 43(2019) 12556-12564
- C. R. Bauwens et al Proceedings of the Combustion Institute 35 (2015) 2059-2066.
- A Ueda et al., Journal of the Energy Institute, 110 (2023) 101335

# Experimental setup



- **Princeton Univ.**

1. Dual chamber
2. Inner cylinder ( $\phi$  82.55 mm x 127 mm)
3. Outer cylinder ( $\phi$  273.05 mm x 304.8 mm)
4.  $V = 0.679\text{L}$ , Up to 60 atm
5. Flame radius 25 mm

- **Hiroshima Univ.**

1. Dual chamber
2. Inner cylinder ( $V = 27.4 \text{ L}$ )
3. Outer cylinder ( $V = 195\text{L}$ )
4. Up to 10 atm
5. Flame radius 110 mm

# Cellular flame images

**Princeton Univ.**

$$P_i = 0.5 \text{ MPa}$$

$$\phi = 0.6$$

$$r = 15.4 \text{ mm}$$

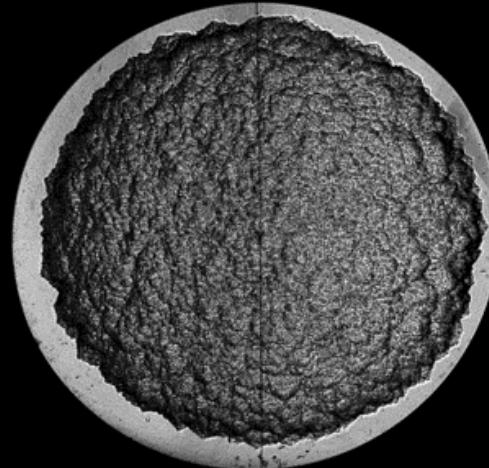


**HIROSHIMA UNIV.**

$$P_i = 0.5 \text{ MPa}$$

$$\phi = 0.6$$

$$r = 100 \text{ mm}$$



- The flame radius  $r = 100 \text{ mm}$ , measured by large dual-chamber in Hiroshima Univ. is much larger than  $r = 15.4 \text{ mm}$  of Princeton Univ.

# Cellular flame images

## Princeton Univ.

$P_i = 0.5 \text{ MPa}$

$\phi = 0.6$

$r = 15.4 \text{ mm}$

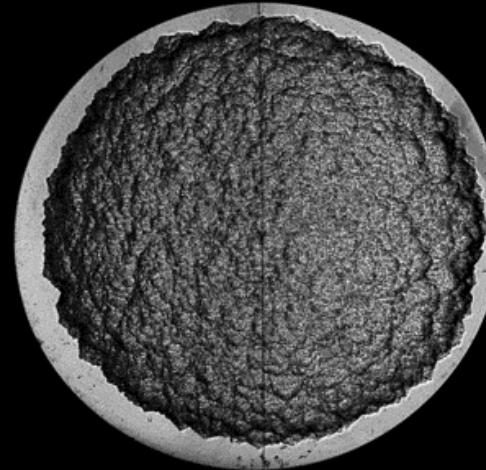


## Hiroshima Univ.

$P_i = 0.5 \text{ MPa}$

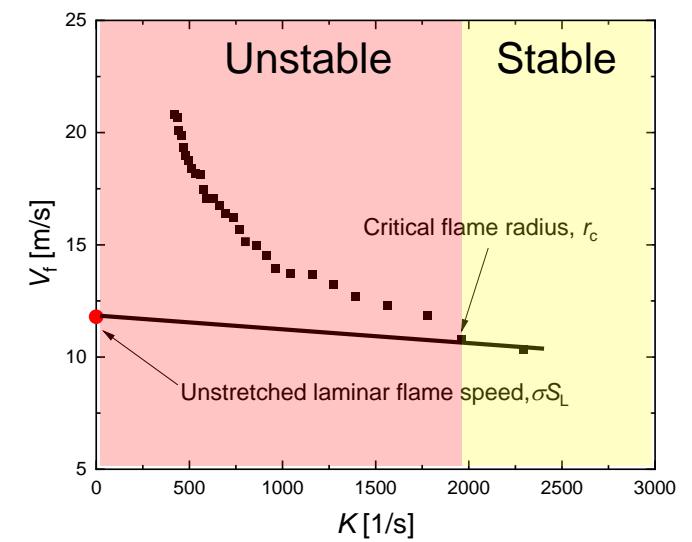
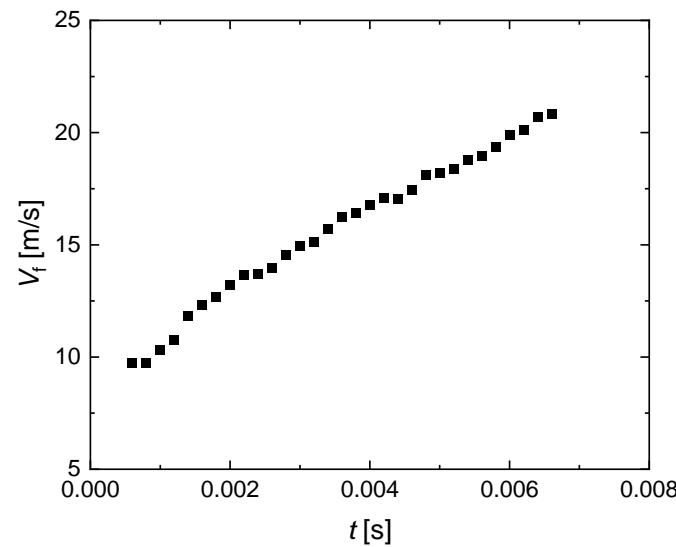
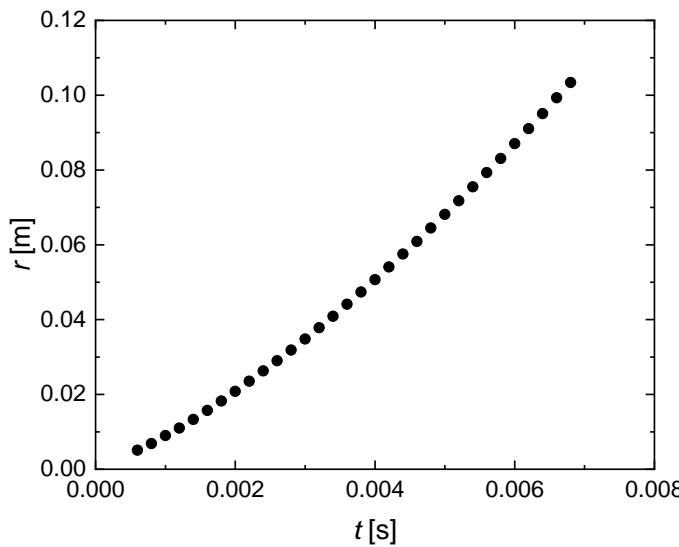
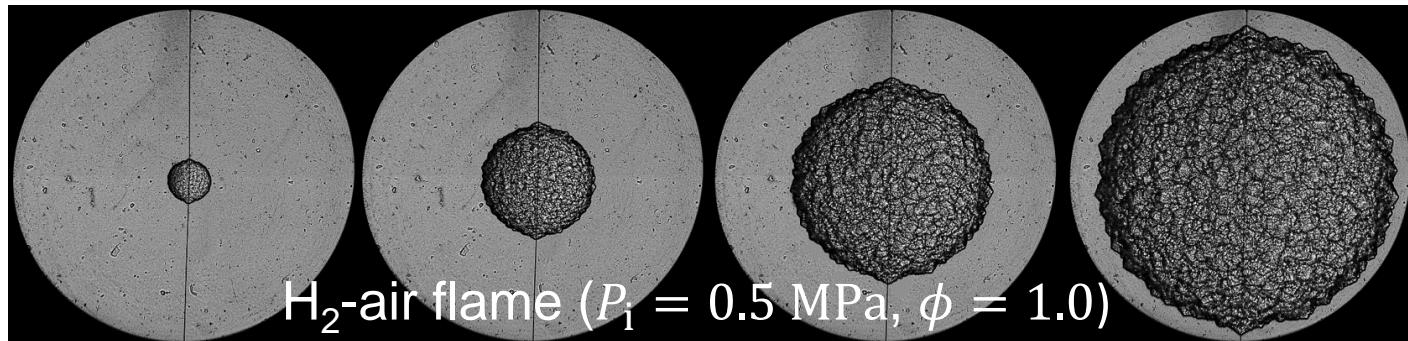
$\phi = 0.6$

$r = 100 \text{ mm}$



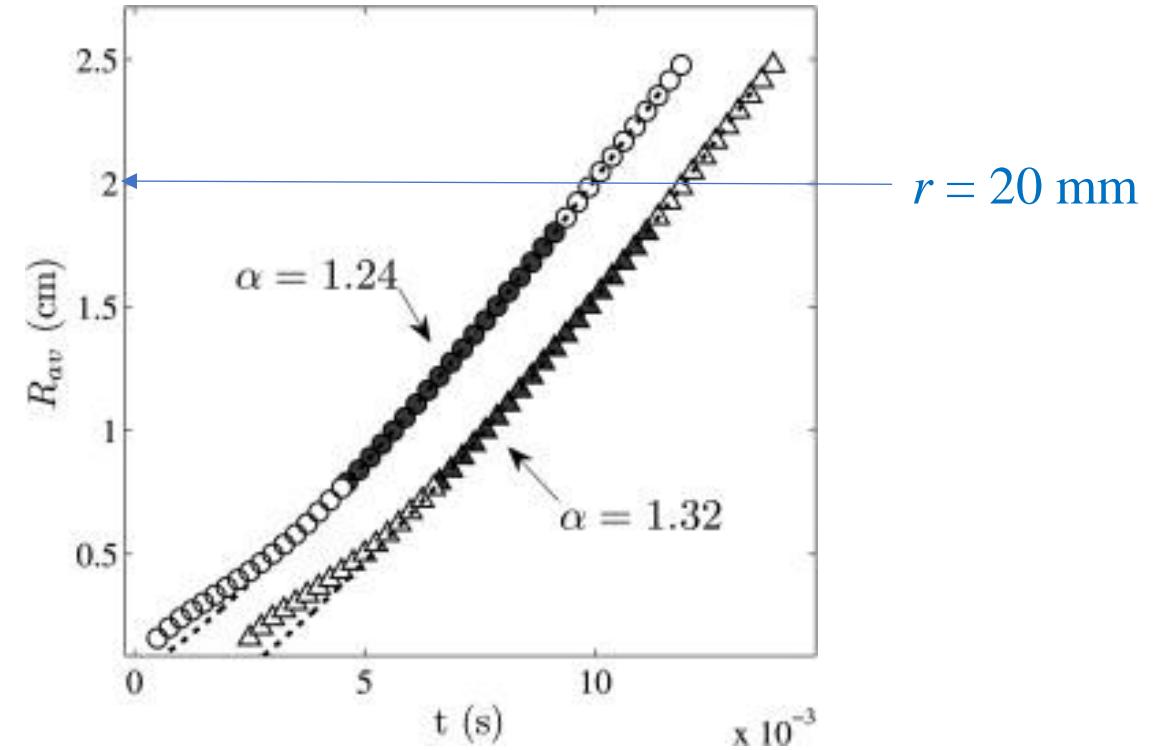
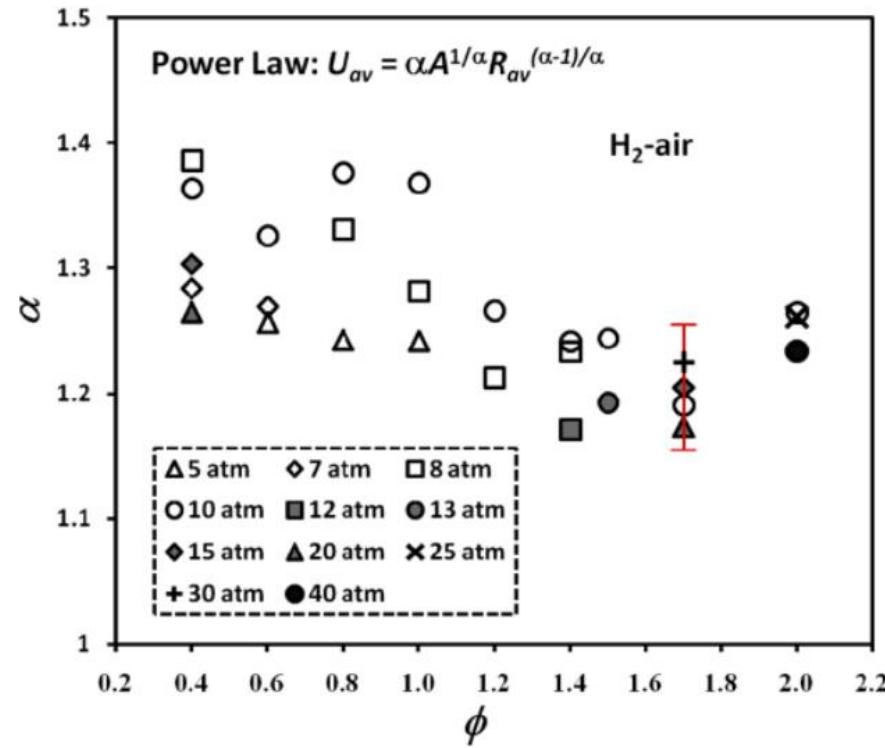
- The flame radius  $r = 100 \text{ mm}$ , measured by large dual-chamber in Hiroshima Univ. is much larger than  $r = 15.4 \text{ mm}$  of Princeton Univ.

# Onset of self-acceleration



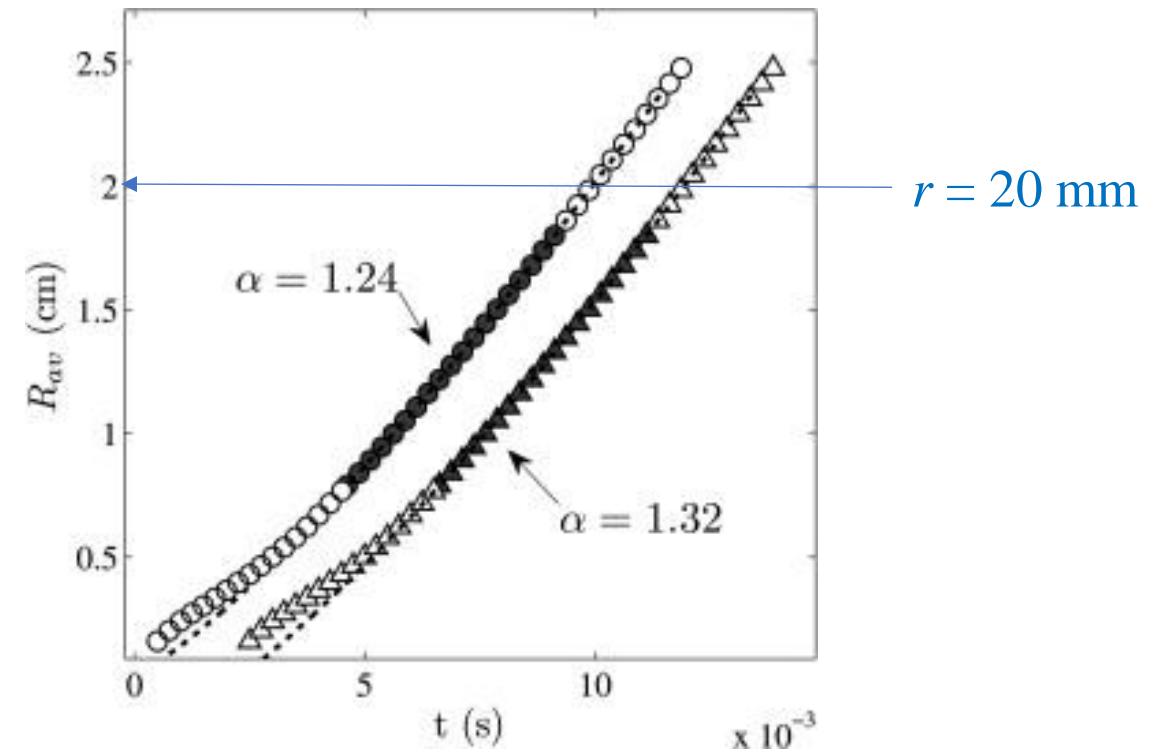
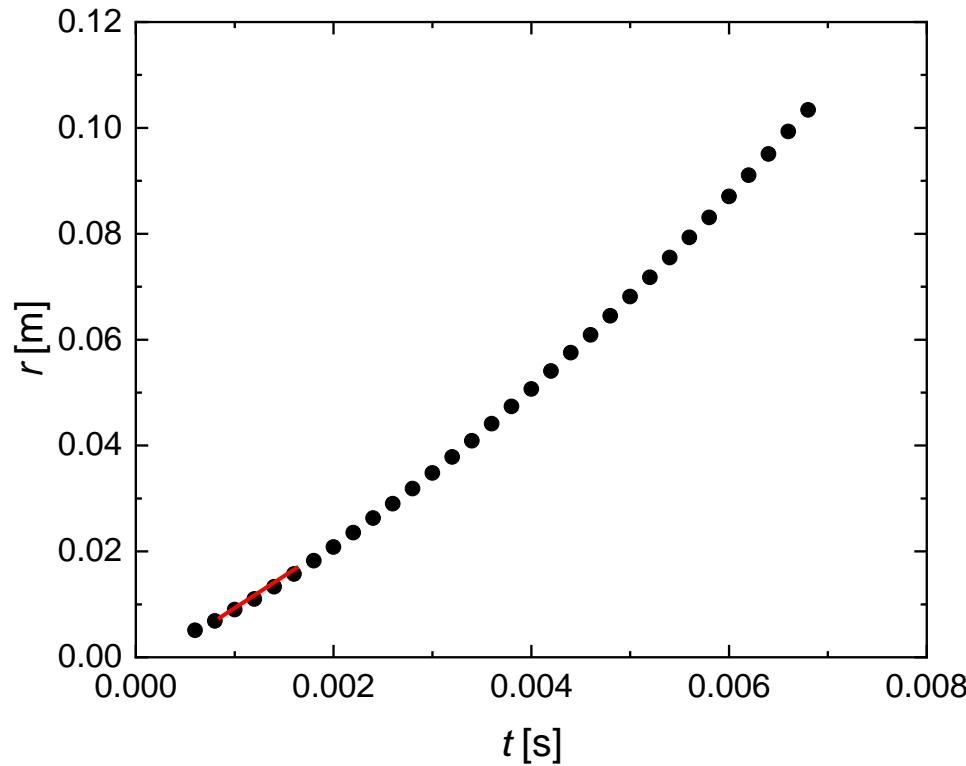
- Spherical expanding flame self-accelerates in stoichiometric  $\text{H}_2\text{-air}$  mixture.

# Acceleration exponent, $r \propto t^\alpha (r \gg r_c)$



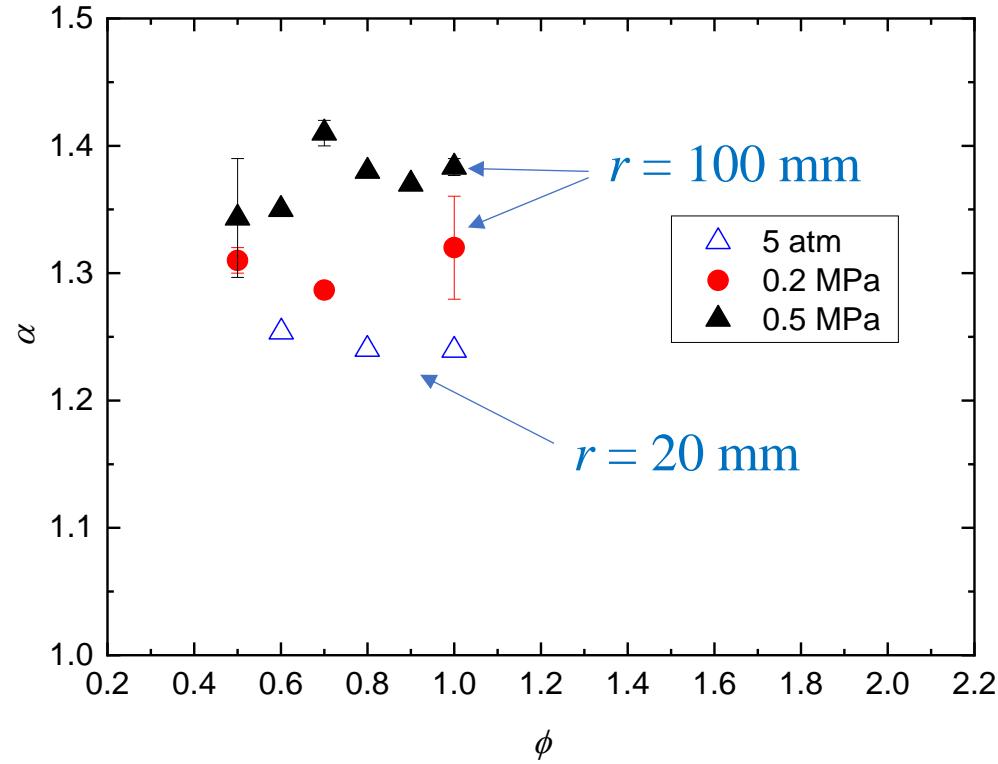
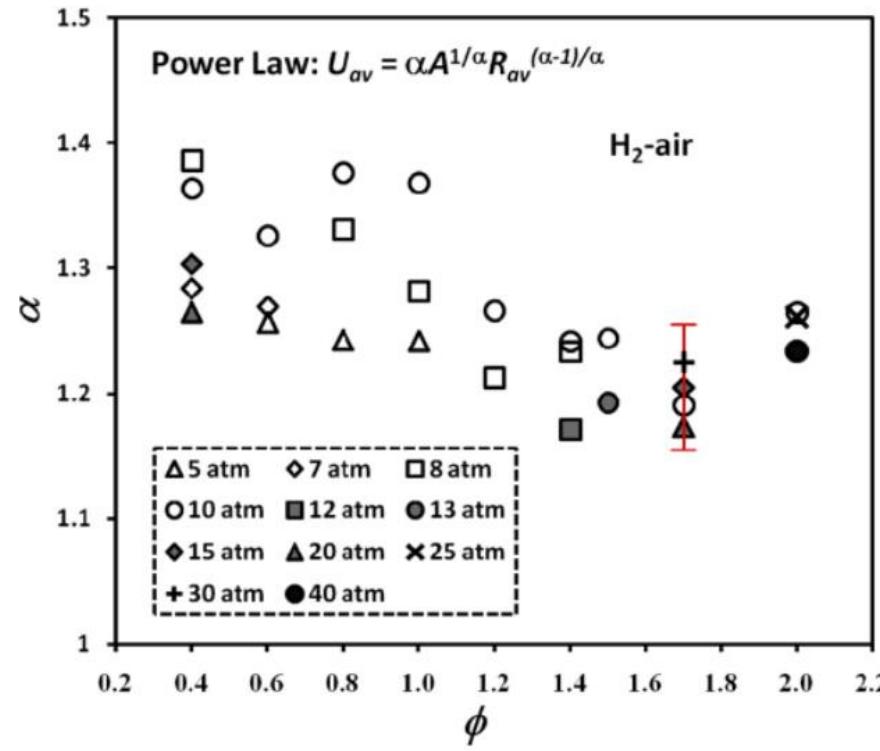
- The increasing tendency of  $\alpha$  with a decrease in  $\phi$ , and  $\alpha$  values increased with initial pressure.
- The  $\alpha$  values seem to depend on the mixture and initial pressure.
- Nevertheless, this result demonstrates that the evaluated values of  $\alpha$  were underestimated, because the evaluation range might be located in the transition regime to self-turbulization.

# Acceleration exponent, $r \propto t^\alpha$ ( $r \gg r_c$ )



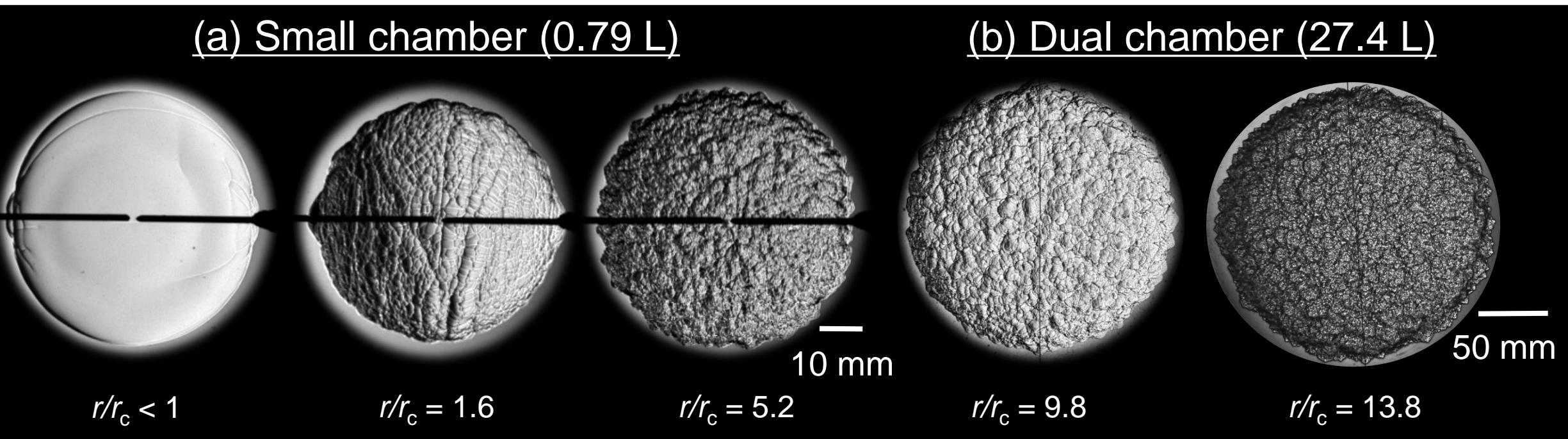
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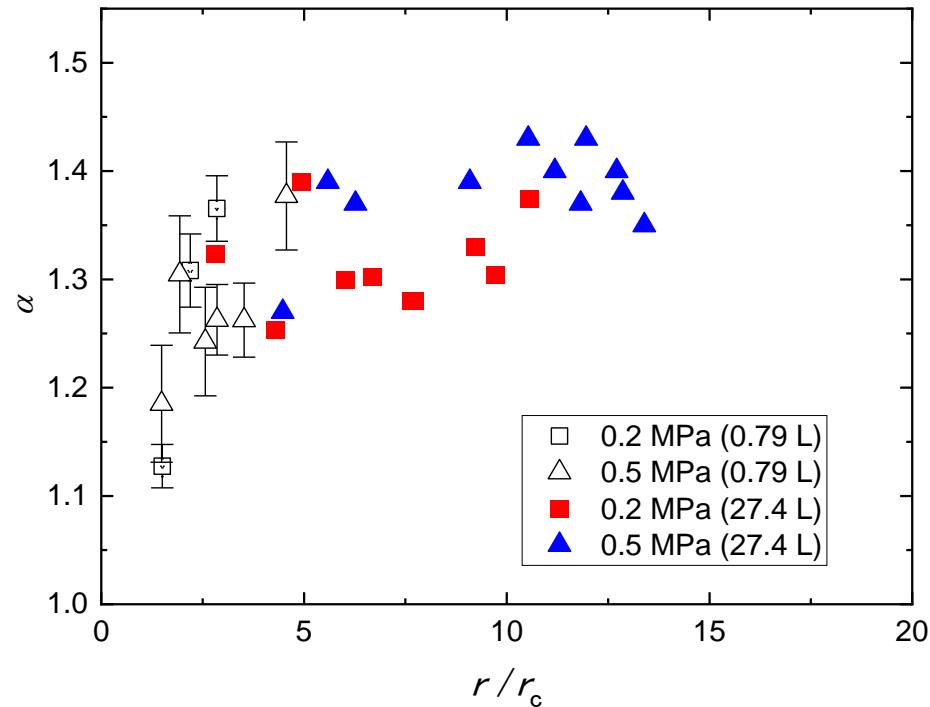
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# Acceleration exponent, $r \propto t^\alpha (r \gg r_c)$



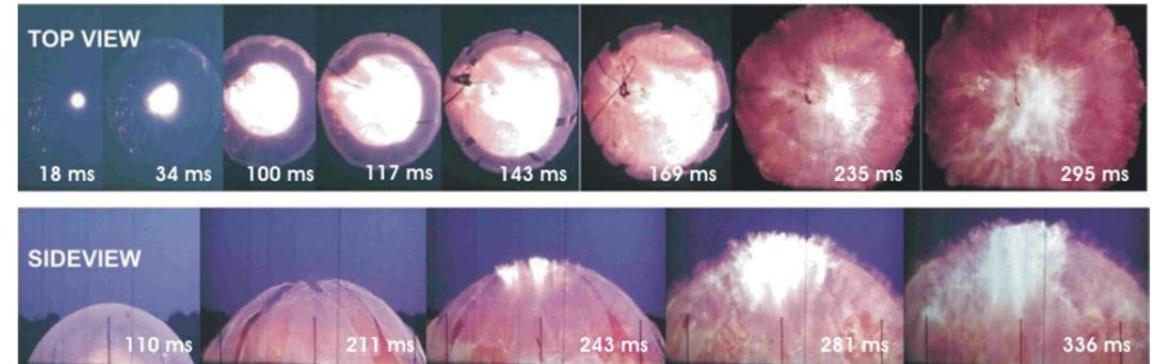
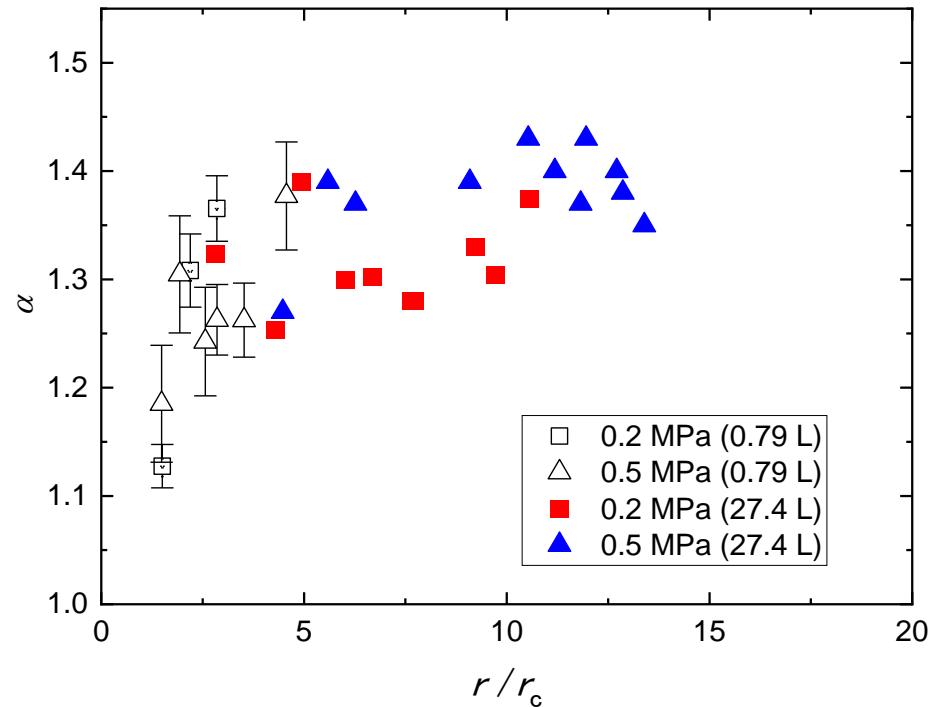
- The  $\alpha$  increased and saturated  $\alpha = 1.4$  with an increase in  $r/r_c$ .
- The transition regime to self-similar propagation has been observed at  $r/r_c > 10$ .
- Self-similarity is observed, in which the value of  $\alpha$  remains nearly constant with further increase in  $r/r_c$ .

# Acceleration exponent, $r \propto t^\alpha (r \gg r_c)$



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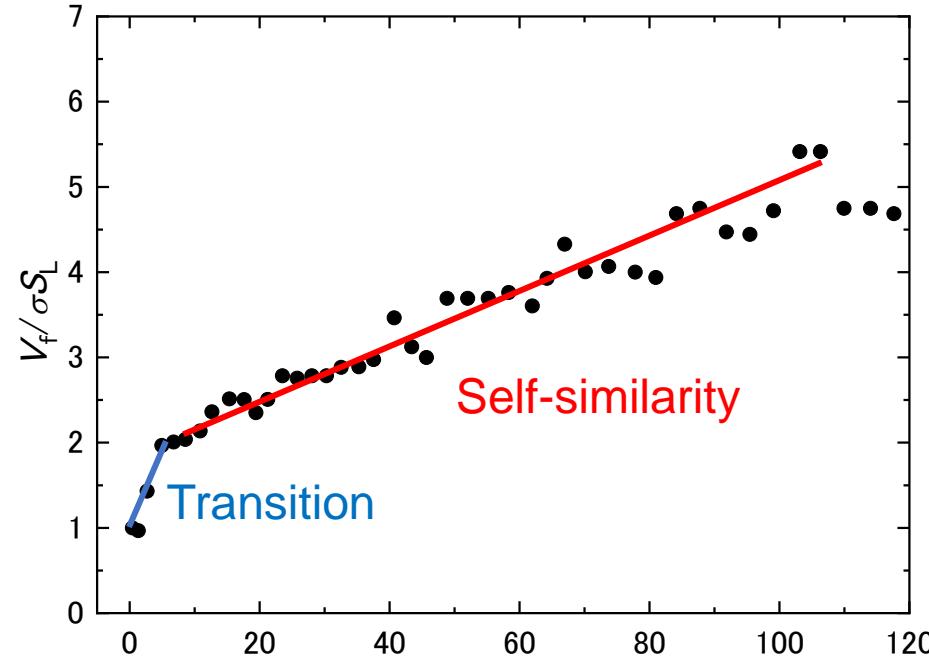
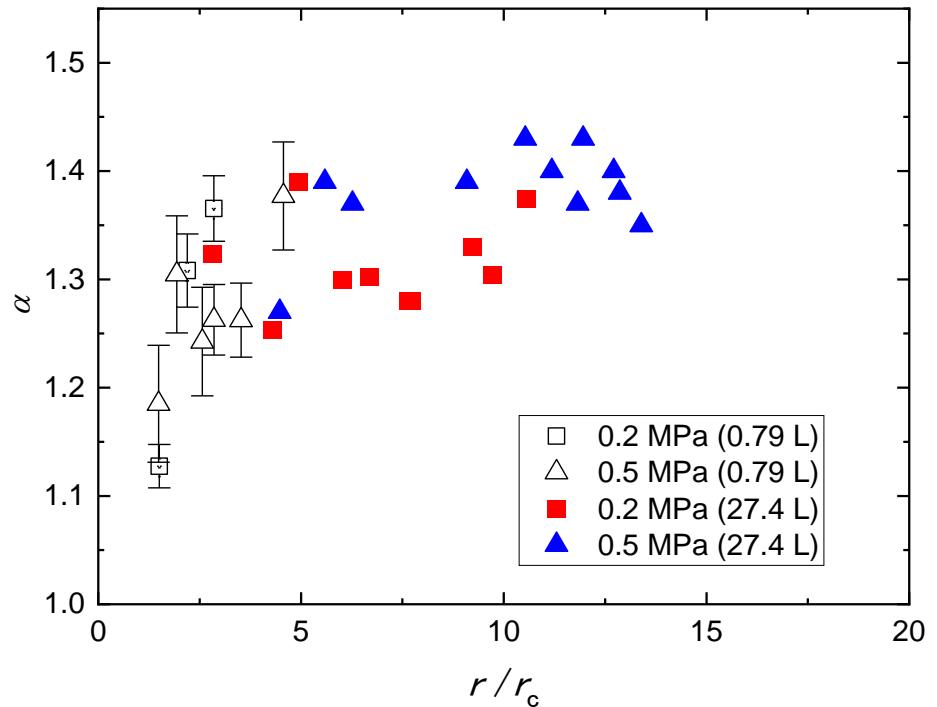
# Acceleration exponent, $r \propto t^\alpha (r \gg r_c)$



Large-scale near stoichiometric hydrogen-air flame in a 20 m diameter hemisphere.

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# Acceleration exponent, $r \propto t^\alpha (r \gg r_c)$

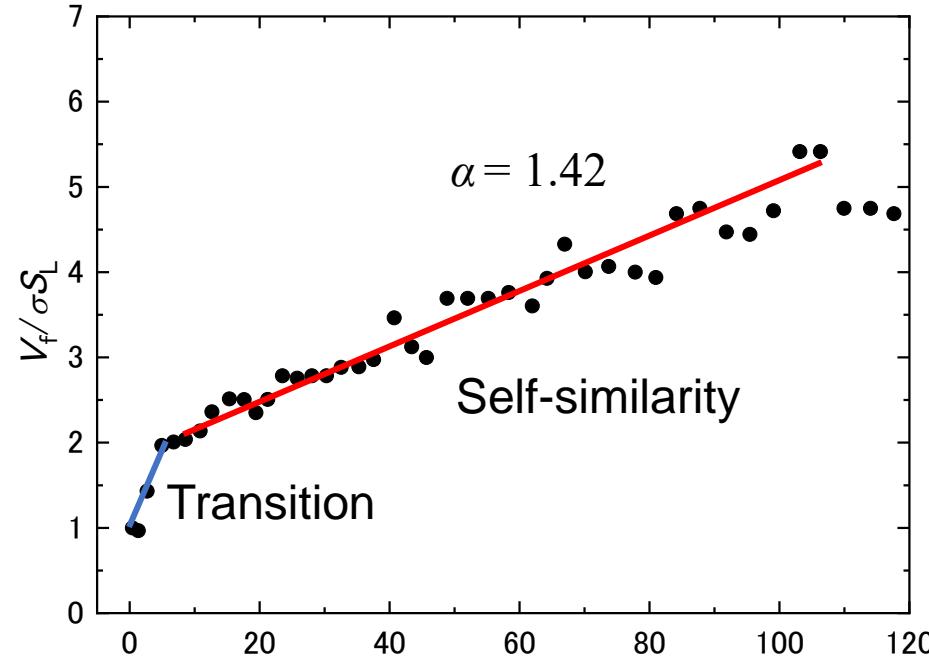
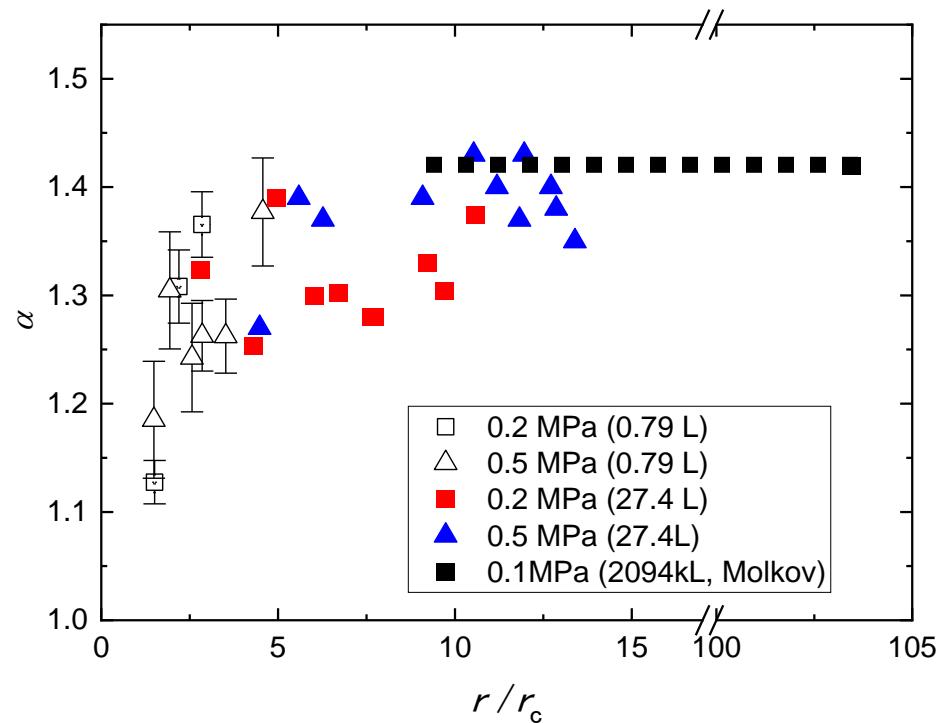


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- The transition regime to self-similar propagation.
- Self-similarity is observed, in which the values of  $V_f / \sigma_S$  increase in  $r/r_c$ .

## Onset of self-similarity

- Molkov,  $r/r_c = 8.6$  ( $r_{cs} = 1$  m)
- Gostinsev,  $r/r_c = 8.6-10.3$  ( $r_{cs} = 1-1.2$  m)
- $r_c = 0.116$  m
- $\alpha = 1.42$

# Acceleration exponent, $r \propto t^\alpha (r \gg r_c)$

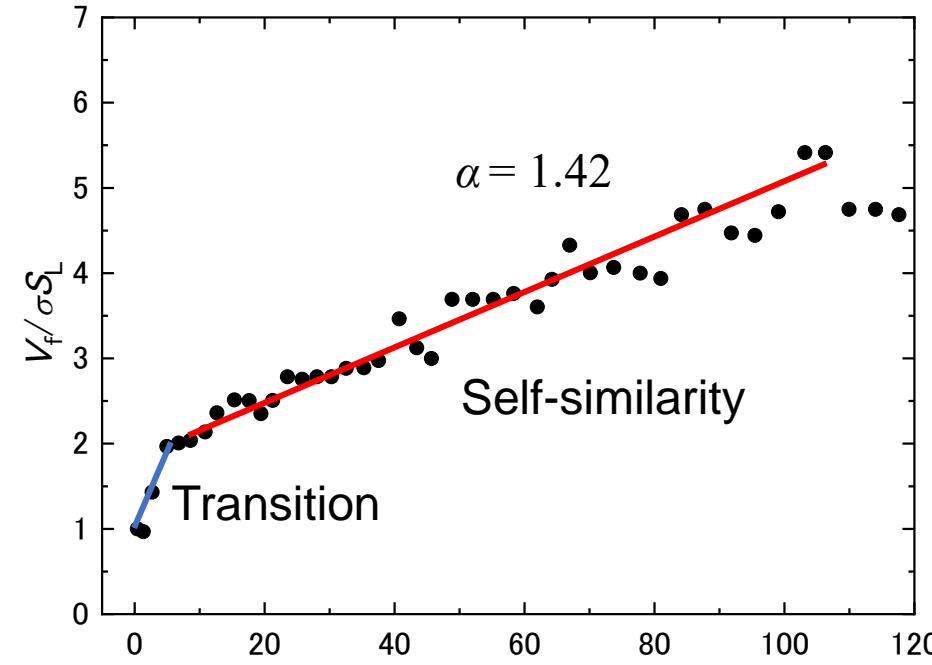
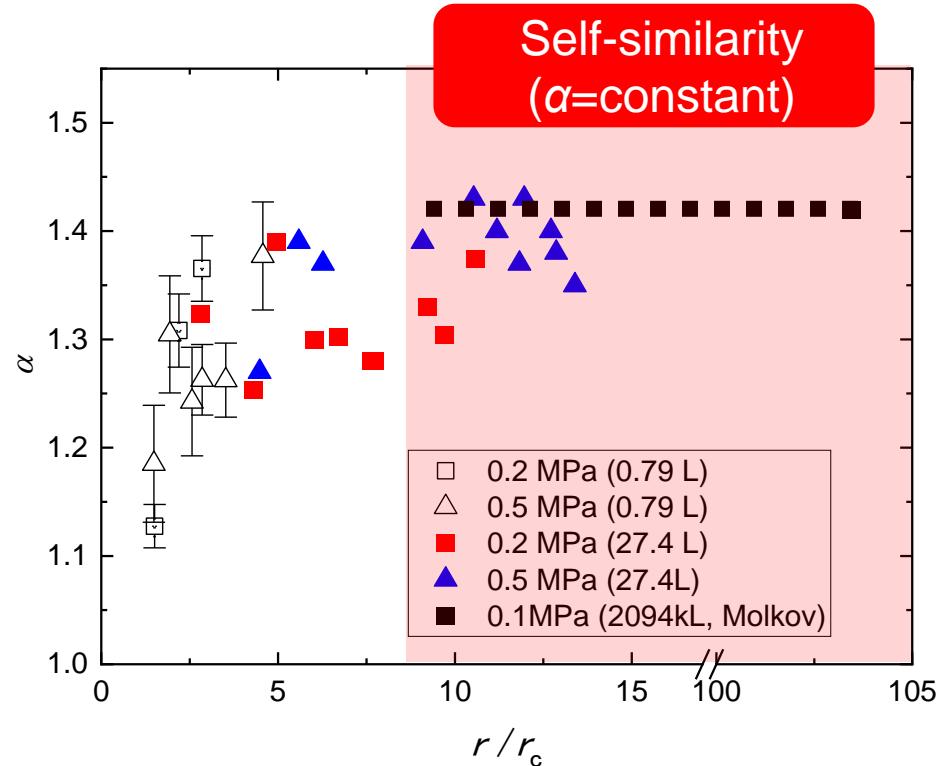


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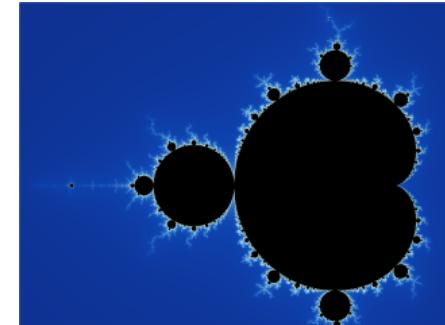
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## Onset of self-similarity

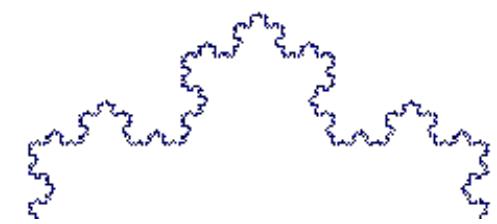
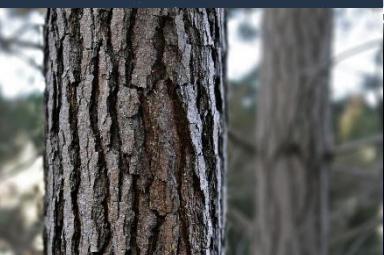
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# Self-acceleration and Self-similarity

## -Fractal pattern in nature-



Mandelbrot set



Koch curve

"Clouds are not spheres, mountains are not cones, coastlines are not circles, and bark is not smooth, nor does lightning travel in a straight line."(Mandelbrot, 1983).

→ **Fractals** are typically **self-similar patterns**, where self-similar means they are "**the same from near as from far**". Fractals may be exactly the same at every scale

# Self-similar formation for Homologous sphere

$$L = \pi r_1$$
$$A = 4\pi r_1^2$$
$$V = \frac{4\pi r_1^3}{3}$$
$$\times \left( \frac{r_2}{r_1} \right)^D$$
$$L = \pi r_2$$
$$A = 4\pi r_2^2$$
$$V = \frac{4\pi r_2^3}{3}$$

$D = 2+d$  : Fractal dimension

# Modification of the flame surface area

Laminar flame propagates spherically

$$r < r_c$$

$$A_L = 4\pi r^2$$

Flame surface area

Onset of turbulence  
 $r_c$ : critical radius

$$r = r_c$$



$$r > r_c$$

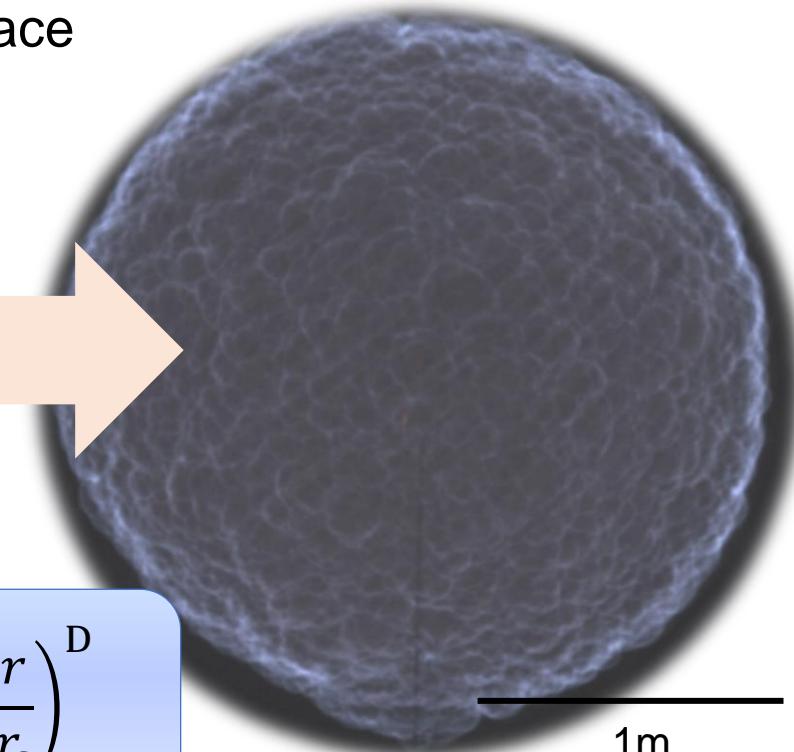
$$A_c = 4\pi r_c^2$$

Similar transformation

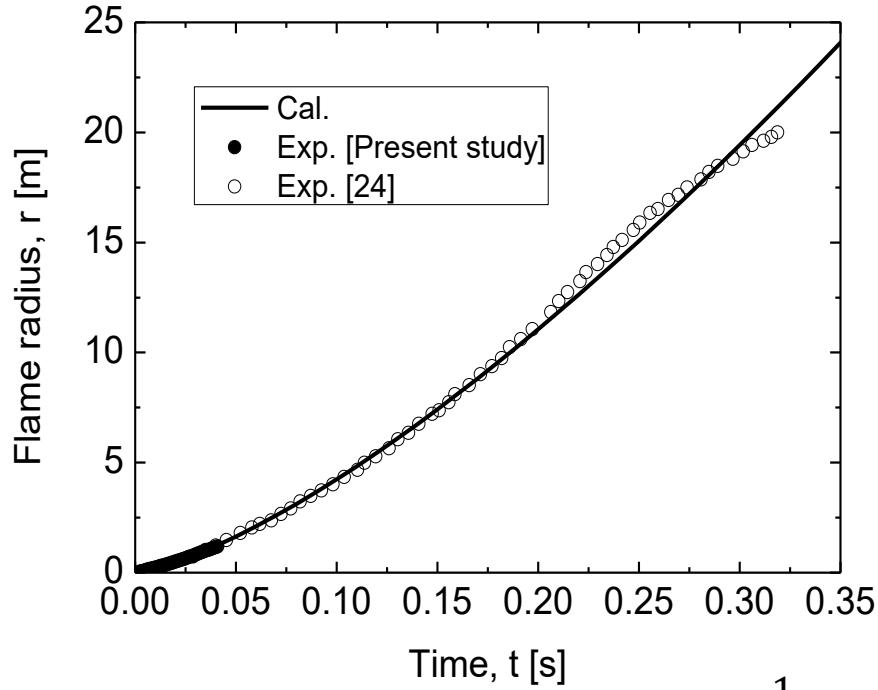
Fractal surface



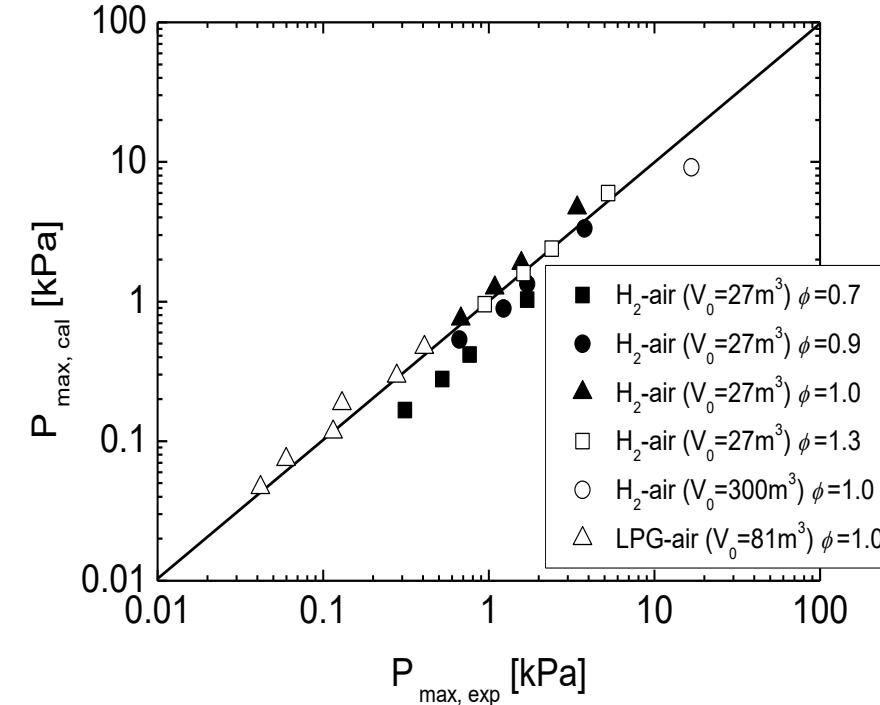
$$A_T = 4\pi r_c^2 \left(\frac{r}{r_c}\right)^D$$



# Estimation models



$$r = \left( \frac{1-d}{r_{\text{cl}}^d} \varepsilon S_L t + d r_{\text{cl}}^{1-d} \right)^{\frac{1}{1-d}}$$



$$p_{\text{max}} = \frac{\rho}{R} (d+2) \varepsilon^2 S_L^2 \frac{r_q^{2d+1}}{r_{\text{cl}}^{2d}}$$

◆ These models were in agreement with large-scale gas explosions.

- W. Kim et al., Int. J. Hydrogen Energy, 40 (2015) 11087-11092,
- V Molkov et al., J. Phys. D: Appl. Phys. 39 (2006) 4366–4376

# -Understanding explosion from industry to space-

